

DE LA RECHERCHE À L'INDUSTRIE



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CHIRPED PULSE AMPLIFICATION FOR SHORT VUV FEL PULSES AT FERMI@ELETTRA

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and

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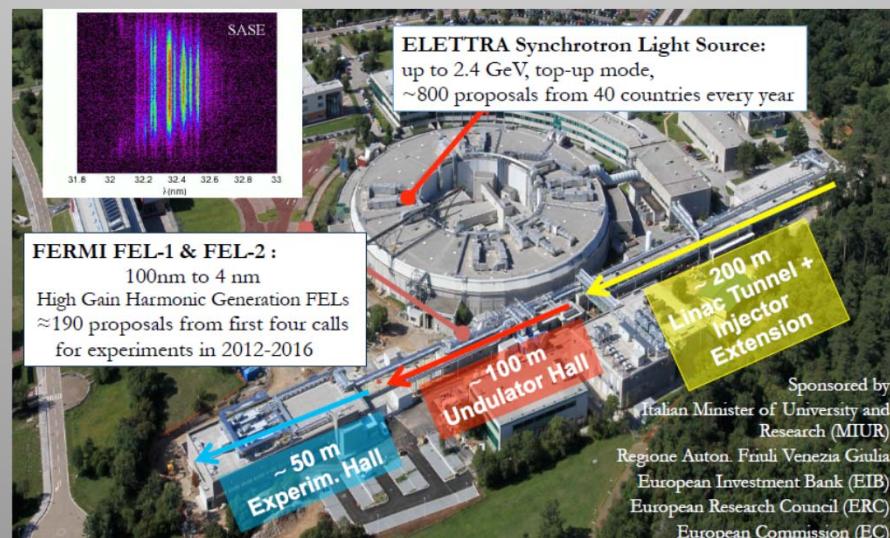
(On behalf of the FERMI CPA Collaboration)

LASERLAB IV Joint research Action

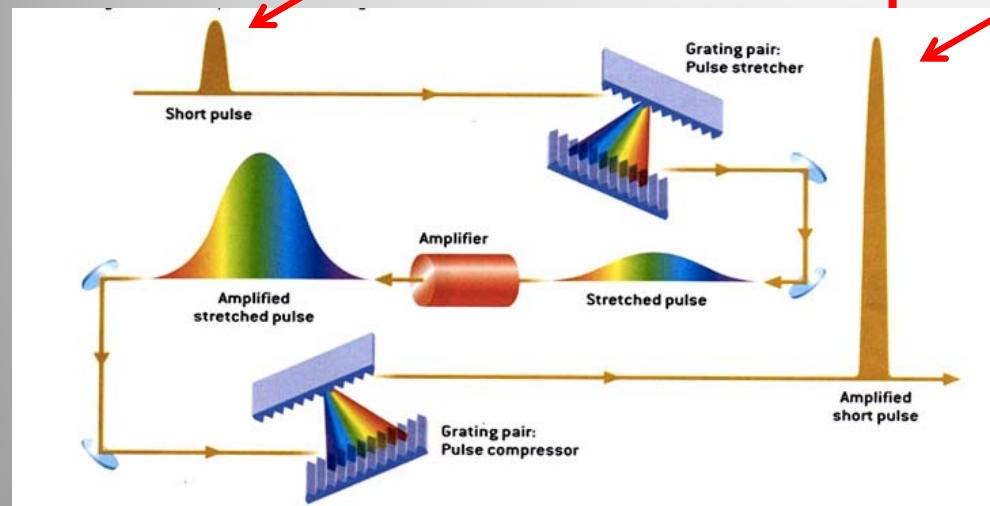
Obj. 3: Advanced coherent XUV and X-Ray sources

- Motivation and theoretical background
- Aim of the first CPA experiment at FERMI :
Scientific, Technical and Operational challenge
 - Electron-beam, seed and FEL preparation
 - The compressor
 - Temporal duration measurement Set-up
- **Results of the first CPA experiment at FERMI and next step**
- Conclusions
- The FERMI CPA collaboration

- Generate sub-femtosecond high peak power XUV pulses
- Fully control the spectro-temporal features of generated radiation
- Extend towards very high photon energies (water window)



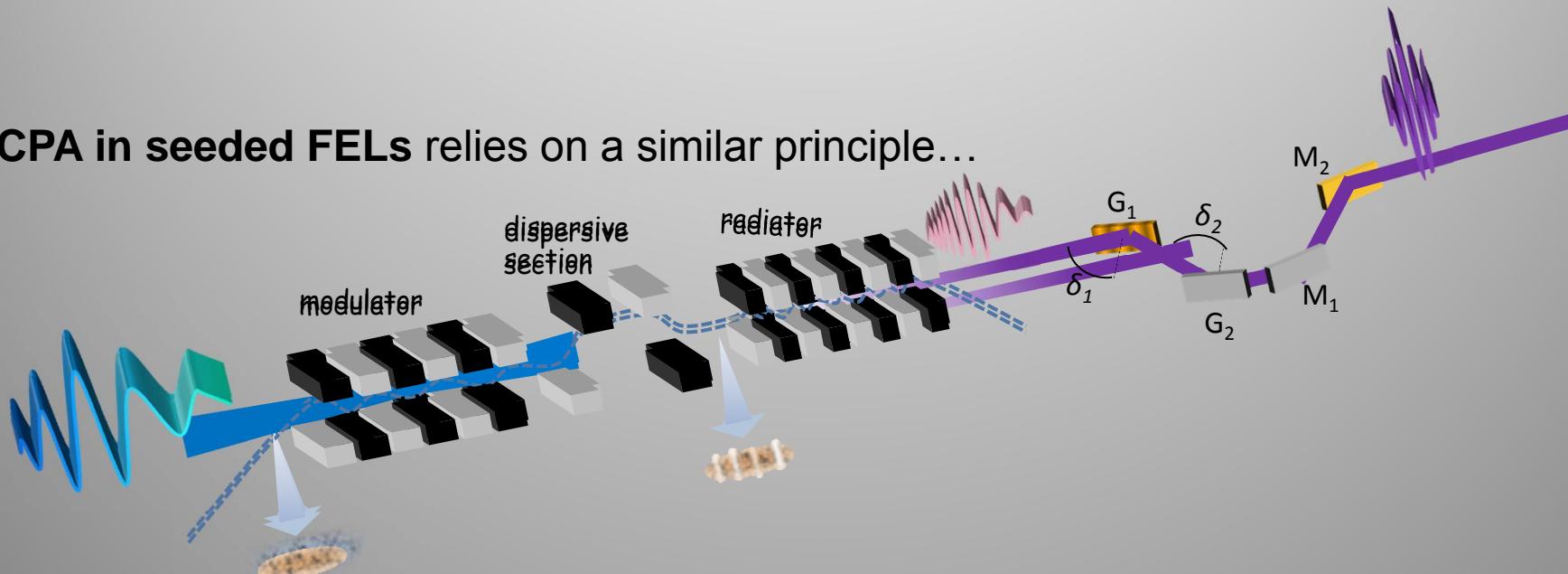
CPA in a solid-state laser....



Similar Bandwidth and input/output duration

Yu, L.H. et al, *Phys Rev. E*, 49 (5), 4480 (1994)
G. Stupakov, SLAC-PUB-14639 (2011)
Ratner, D. et al.. *Phys. Rev. STAB* 15, 030702 (2012).

CPA in seeded FELs relies on a similar principle...



In general, the electric field of a chirped pulse reads

$$\hat{E}(\omega) \sim \exp\left(-\frac{\omega^2}{2\sigma_\omega^2}\right) \exp\left(-\frac{1}{4}i\beta\omega^2\right),$$

spectral width

group delay dispersion (GDD)

In the time domain

$$E(t) \sim \exp\left(-\frac{t^2}{2\sigma_t^2}\right) \exp\left(-i\Gamma t^2\right),$$

temporal width

chirp coefficient

$$\begin{cases} \Gamma = \frac{\beta}{4/\sigma_\omega^4 + \beta^2} \\ \sigma_t^2 = \frac{1}{\sigma_\omega^2} + \frac{1}{4}\beta^2\sigma_\omega^2 \end{cases}$$

$$\beta \ll 2/\sigma_\omega^2$$

small-chirp regime

$$\begin{cases} \Gamma \approx 0 \\ \sigma_t \approx \frac{1}{\sigma_\omega} \quad (\text{transform limit}) \end{cases}$$

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temporal width

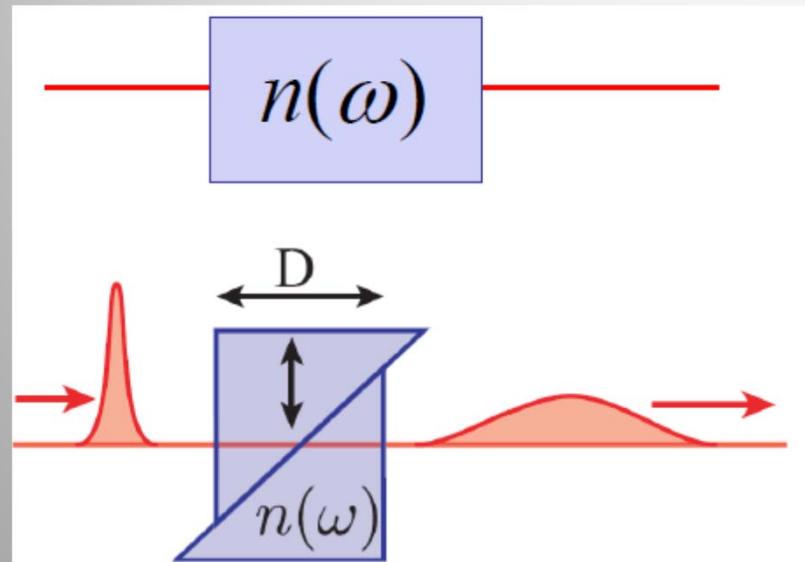
chirp coefficient

$$\begin{cases} \Gamma = \frac{\beta}{4/\sigma_\omega^4 + \beta^2} \\ \sigma_t^2 = \frac{1}{\sigma_\omega^2} + \frac{1}{4}\beta^2\sigma_\omega^2 \end{cases}$$

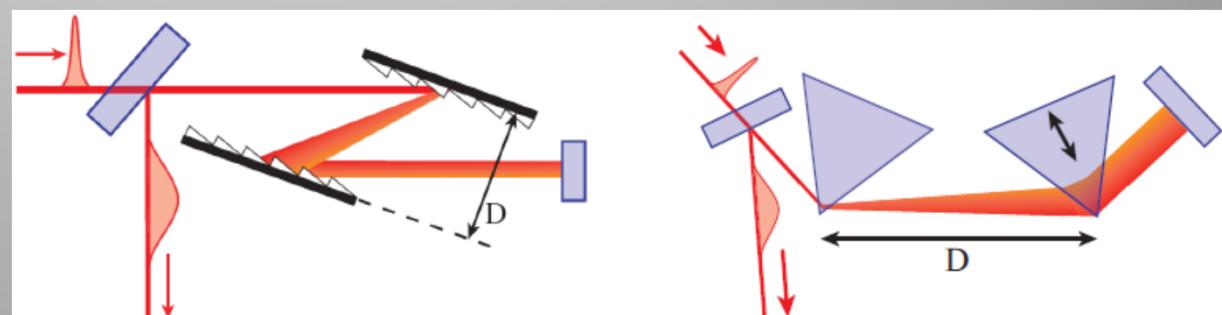
$\beta \gg 2/\sigma_\omega^2$

strong-chirp regime

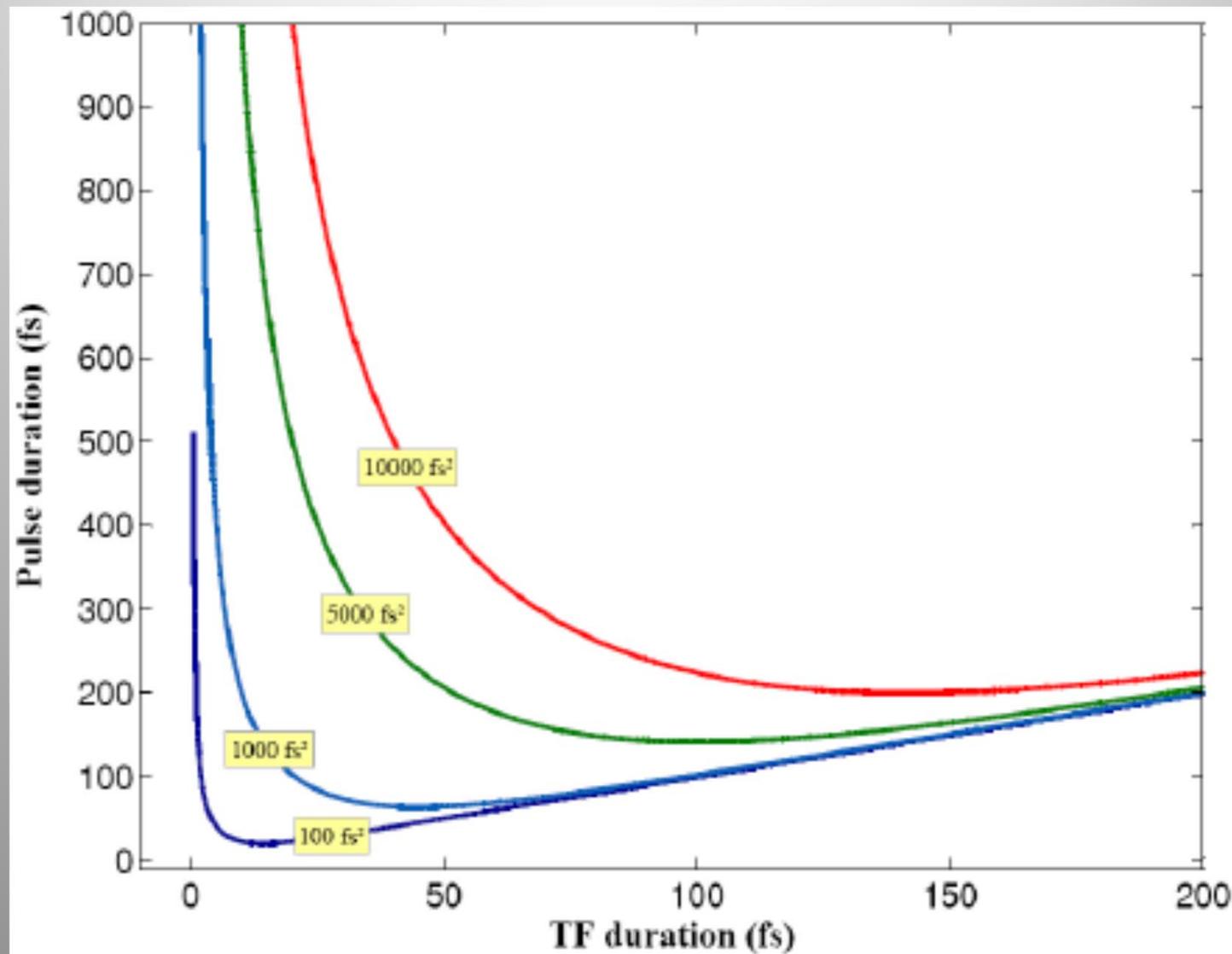
$$\begin{cases} \Gamma \approx \frac{1}{\beta} \\ \sigma_t \approx \frac{1}{2}\beta \sigma_\omega \end{cases}$$



High Dispersive Materials



Prisms or Diffraction Gratings pairs



Strong-chirp regime

$$\begin{cases} \sigma_t \approx \frac{1}{2} \beta \sigma_\omega \\ \Gamma \approx \frac{1}{\beta} \end{cases}$$

Strongly-chirped seed \longrightarrow Strongly-chirped FEL

Moreover :

$$(\sigma_t)_{FEL}^{noCPA} = \frac{(\sigma_t)_{seed}}{n^\alpha} \quad n = \frac{\lambda_{FEL}}{\lambda_{seed}} \approx 10 \quad \text{harmonic number}$$

$$\alpha \approx 1/3$$

(for “moderate” FEL saturation: Stupakov’s law)

$$\Gamma_{FEL} = n \Gamma_{seed} \quad \longrightarrow \quad \beta_{FEL} \approx \beta_{seed} / n$$

Combining previous relations...

$$(\sigma_\omega)_{FEL} = n^{1-\alpha} (\sigma_\omega)_{seed}$$

FEL bandwidth larger than seed bandwidth!



seed pulse duration (***no chirp***)

$$(\sigma_t)_{FEL}^{CPA} = \frac{(\sigma_t)_{seed}^{nochirp}}{n^{1-\alpha}}$$

FEL pulse duration after compression

CPA

$$(\sigma_t)_{FEL}^{CPA} = \frac{(\sigma_t)_{seed}^{nochirp}}{n^{1-\alpha}}$$

no CPA

$$(\sigma_t)_{FEL}^{noCPA} = \frac{(\sigma_t)_{seed}}{n^\alpha}$$

Typically $\alpha \approx 1/3$

$$(\sigma_t)_{FEL}^{CPA} \approx -\frac{1}{n^{2/3}}$$

$$(\sigma_t)_{FEL}^{noCPA} \approx -\frac{1}{n^{1/3}}$$



Demonstrate the feasibility of the CPA scheme in a seeded FEL



Check the validity of the scaling law

$$(\sigma_t)_{FEL}^{CPA} = \frac{(\sigma_t)_{seed}^{nochirp}}{n^{1-\alpha}}$$

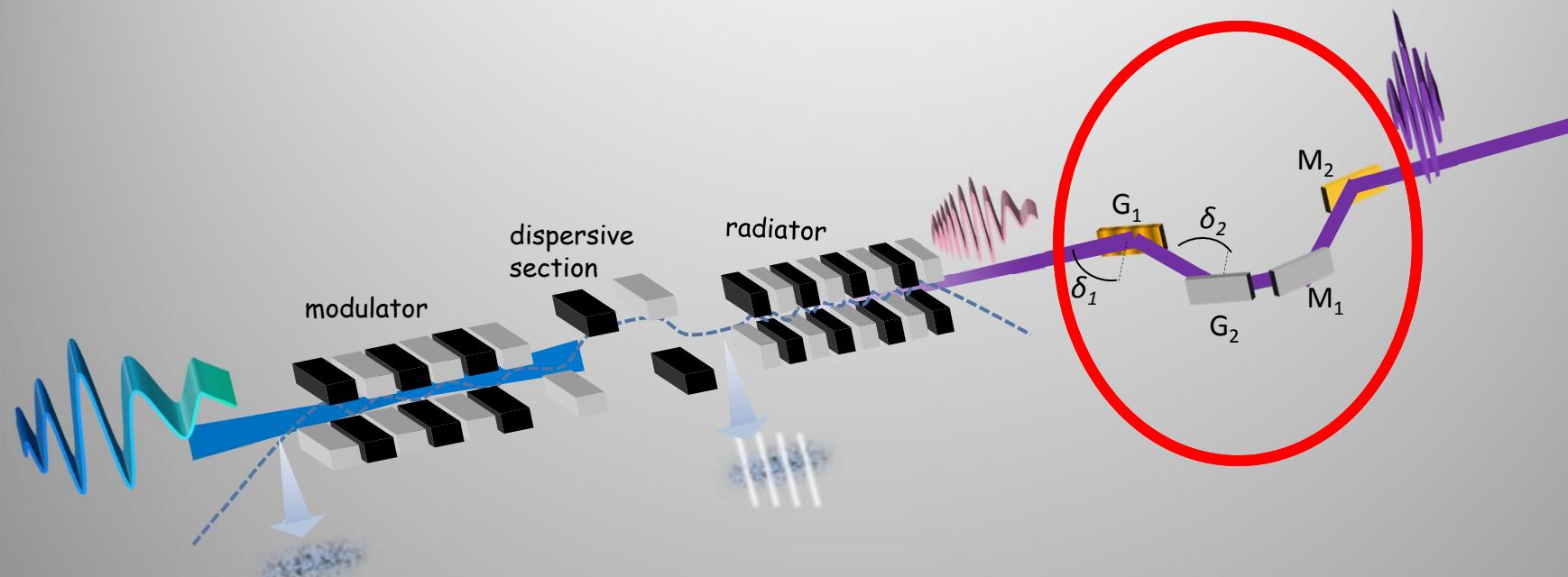


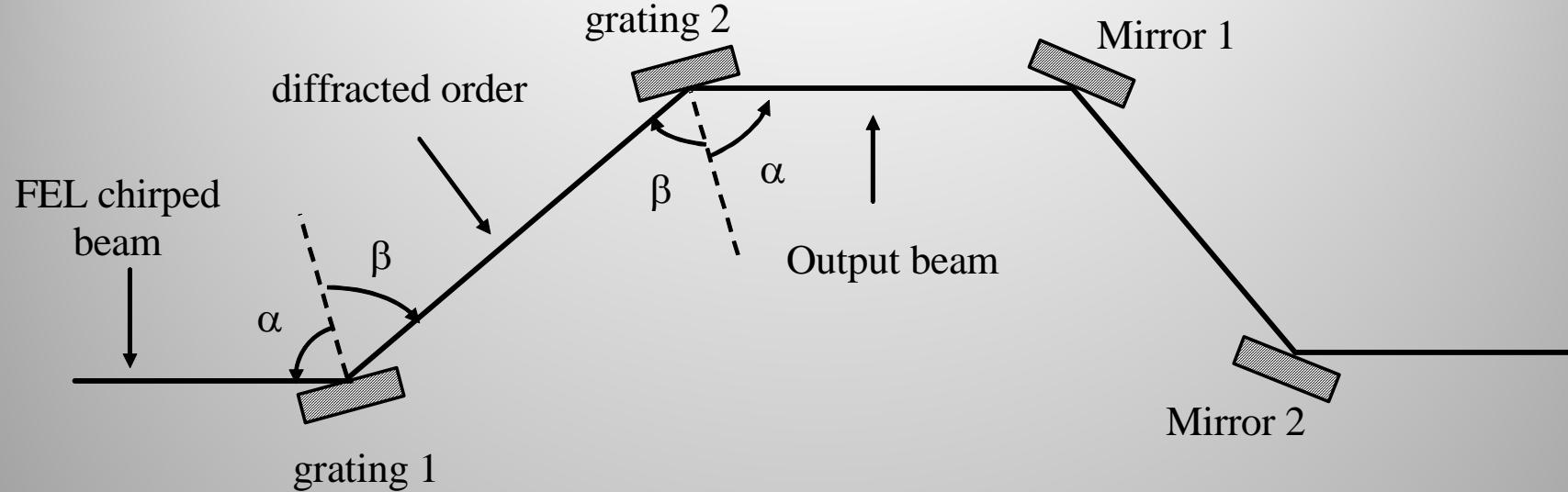
Optimize every single stage of the experimental Setup

- Electron and Laser seed pulse
- The Compressor
- Pulse Duration Measurement Techniques

The compressor

The compressor



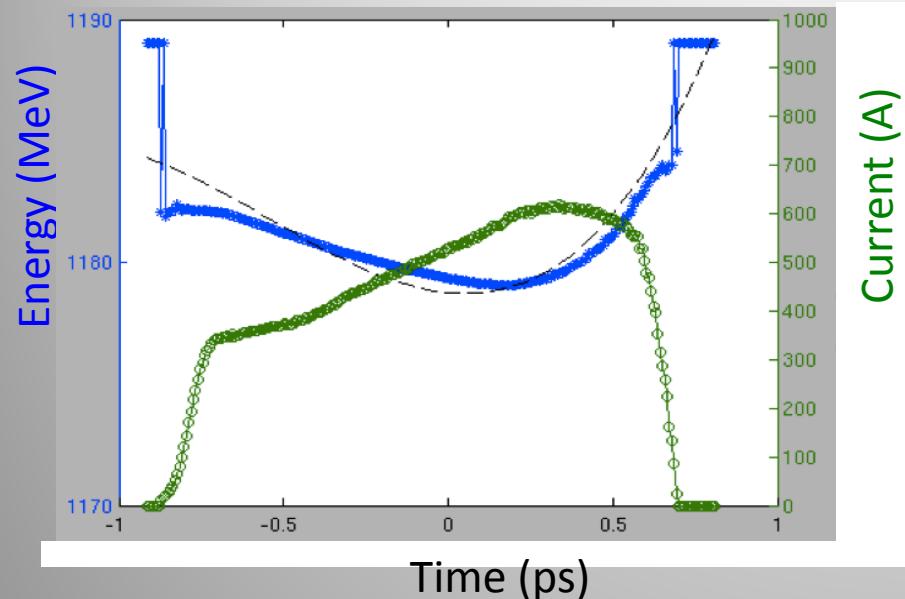


Spectral range	20-40 nm
Spectral bandwidth	0.02-0.015 nm
Temporal compensation	100-400 fs
Dimensions	<1.5 m
Measured efficiency	$\approx 5\%$

Electron-beam, seed and FEL preparation

Control the Energy electrons Chirp

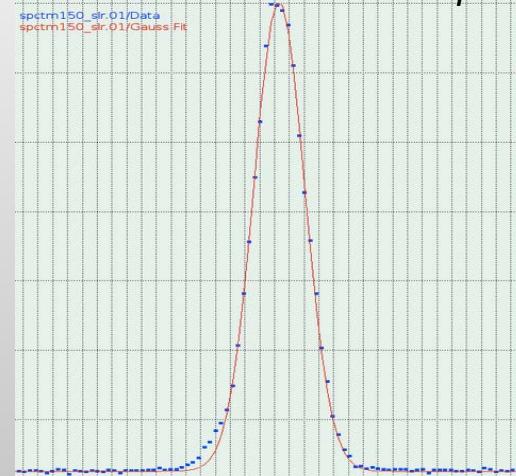
$$\lambda_{\text{res}} = \frac{\lambda_0}{2n\gamma^2} \left(1 + \frac{K^2}{2} \right)$$



Electron-beam parameters	
Bunch duration	1.3 ps
Peak current	400-600 A
Energy	1.18 GeV
Energy spread (slice)	150-200 KeV
Emittance (projected)	2.7/1.1 mm mrad
Quadratic energy chirp	few Mev/ps ²

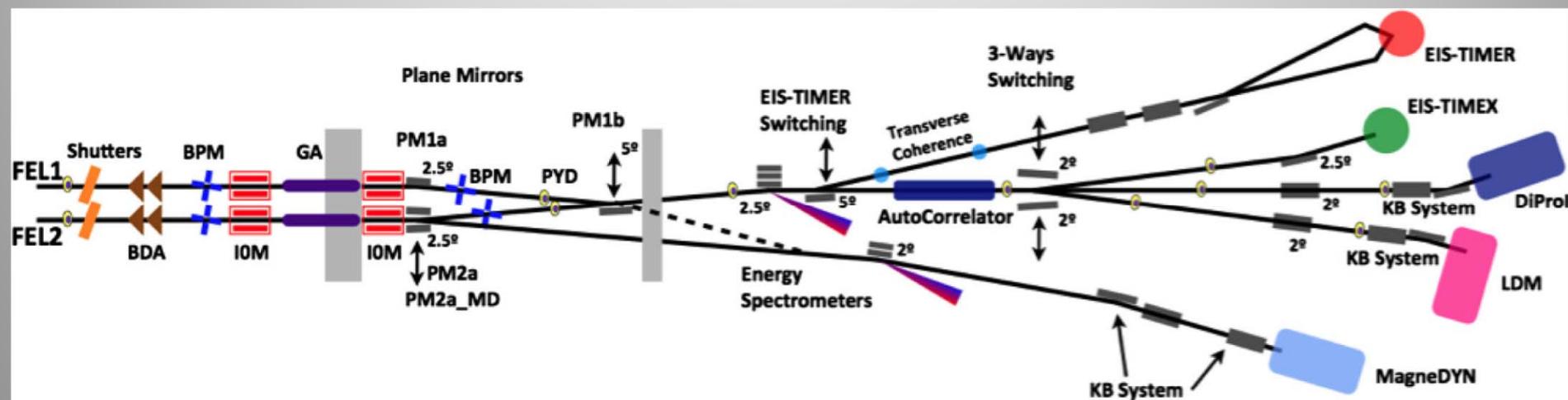
Control the Seed stretching

Non-stretched seed pulse



Seed parameters	
Wavelength	261 nm
Pulse duration (before stretching)	170 fs (FWHM)
Pulse duration (after stretching)	290 fs (FWHM)
Pulse bandwidth	0.7 nm
Power	100-200 MW

Measurement of FEL pulse duration: method and analysis



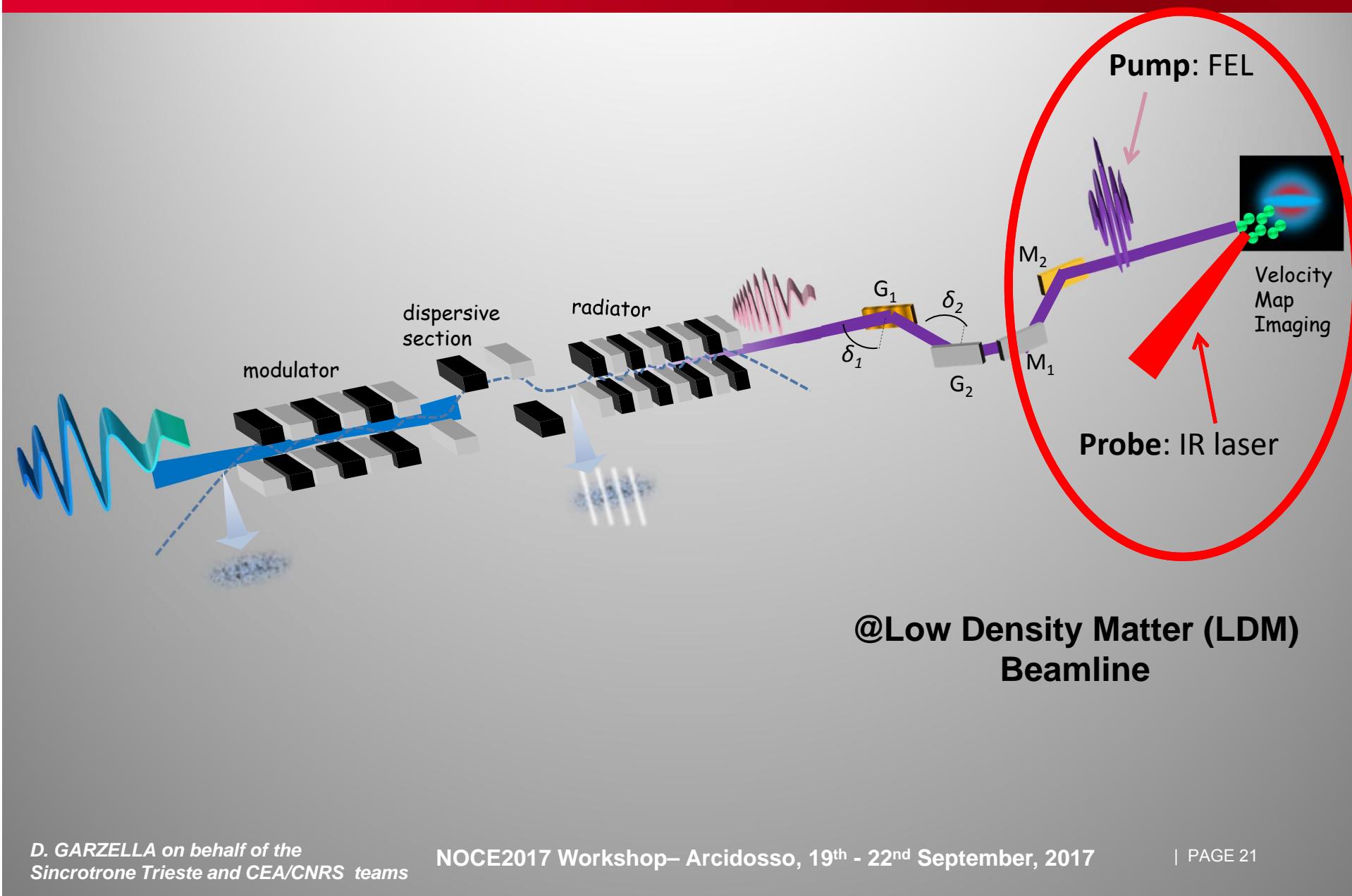
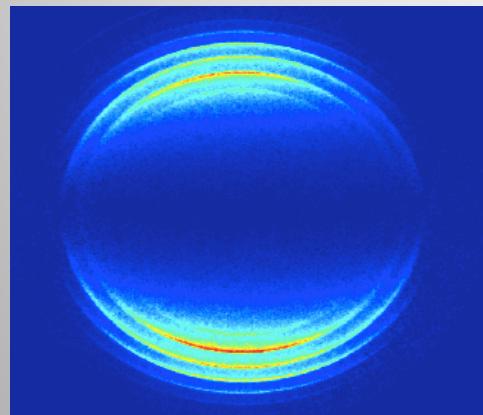
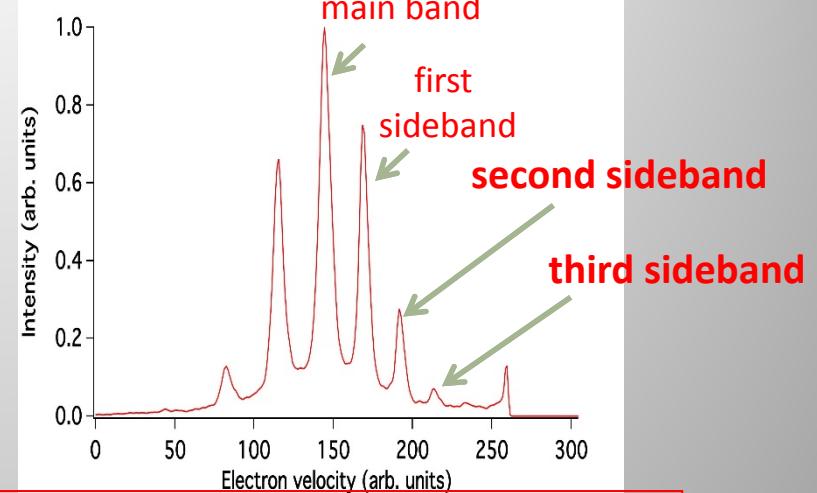
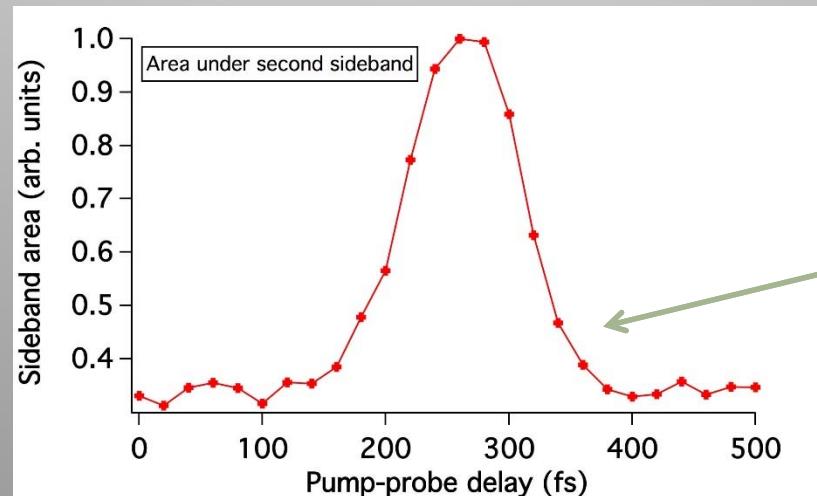


Photo-electron distribution**Energy spectrum****Cross-correlation curve by varying the pump-probe delay**

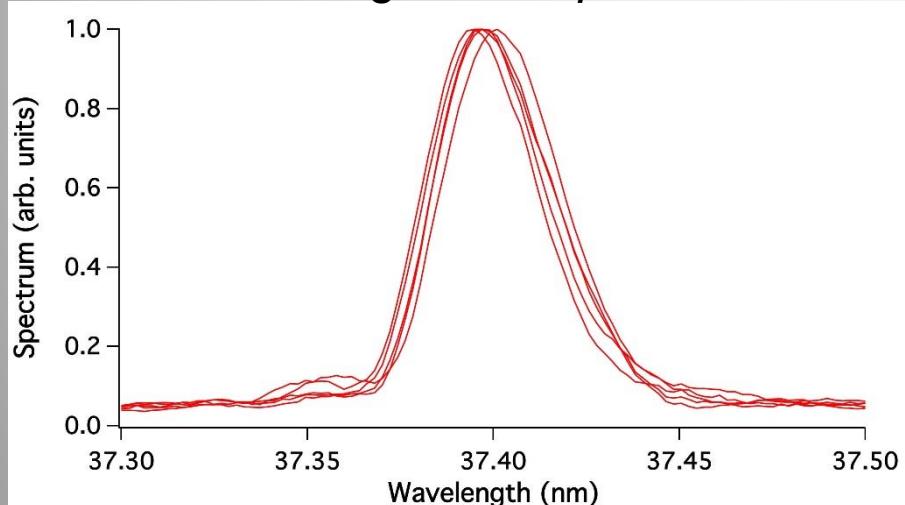
*FEL pulse duration
obtained by
deconvolving the
duration of the
Laser seed pulse*

Results of the CPA experiment

The initial FEL pulse (no seed stretching, no FEL compression)

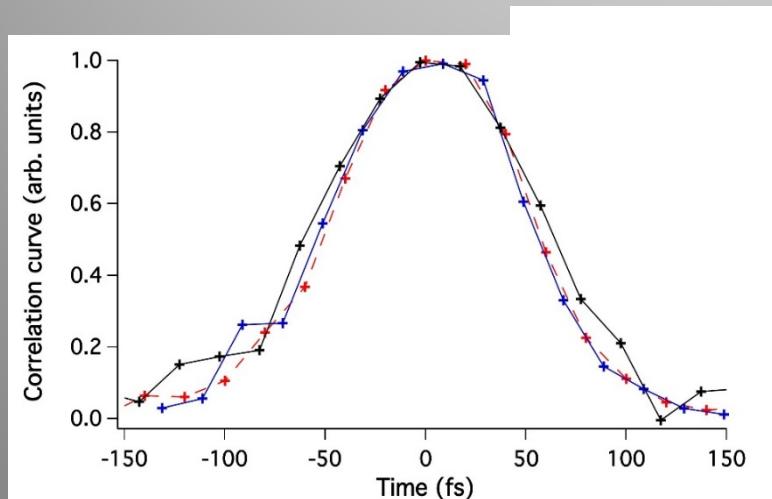


Single-shot spectra



Average spectral width: 0.0384 nm

Cross-correlation curves



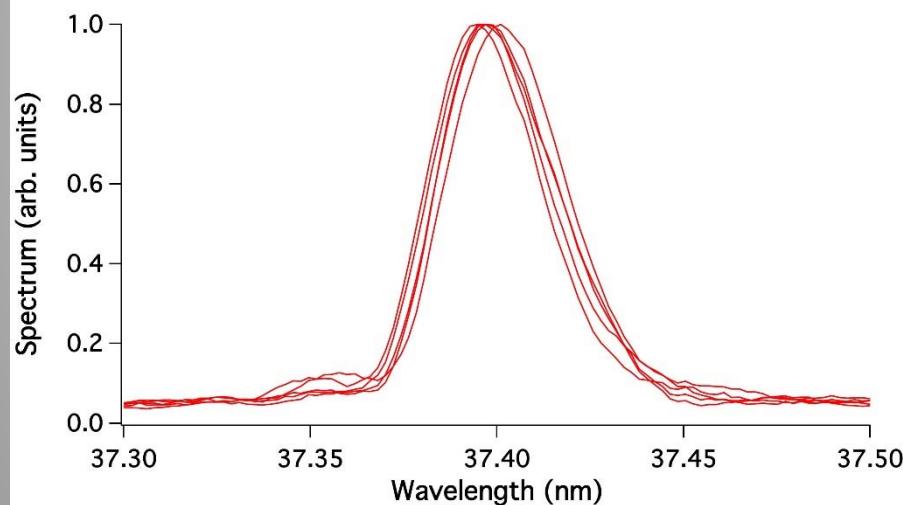
FEL pulse duration: 91 fs

Good agreement with expectations:
 $(\sigma_t)_{\text{seed}}/n^{1/3} = 89 \text{ fs}$

Time-bandwidth product
factor 1.7 above transform limit

The stretched FEL pulse (no compression)

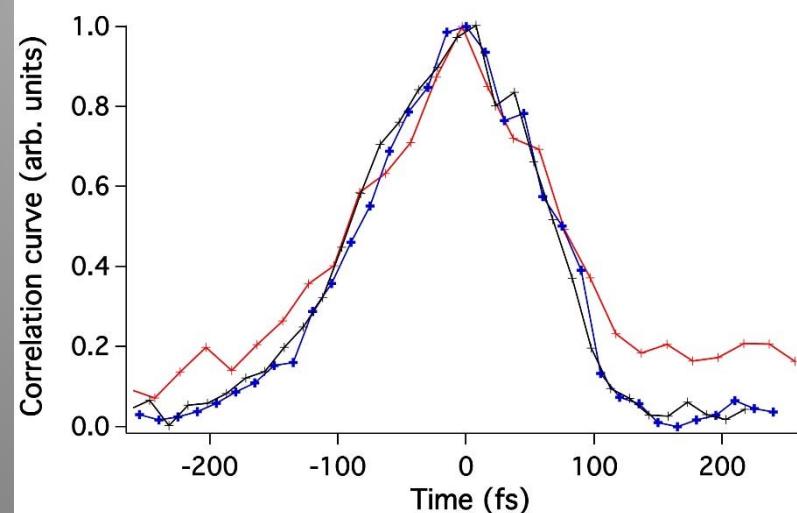
Single-shot spectra



Spectral width: 0.0446 nm

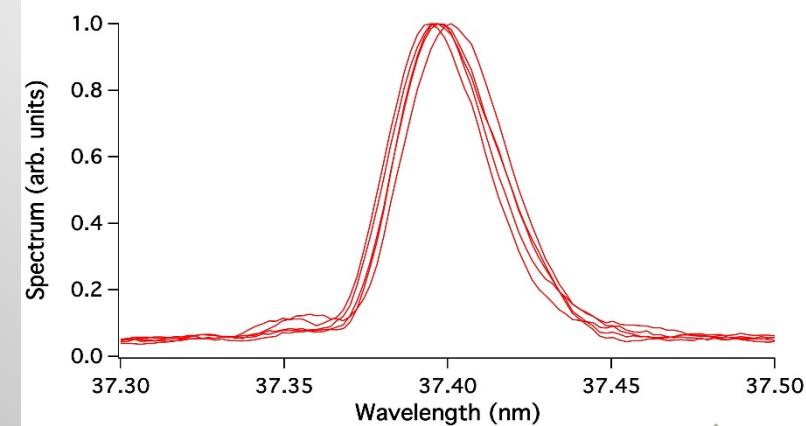
Fairly good agreement with the scaling law predicted by the CPA theory

Cross-correlation curves



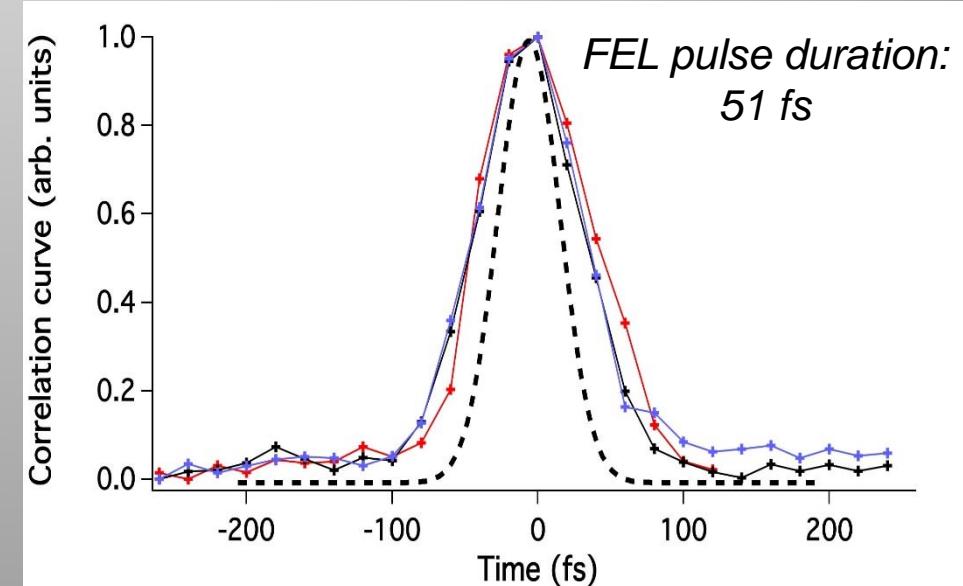
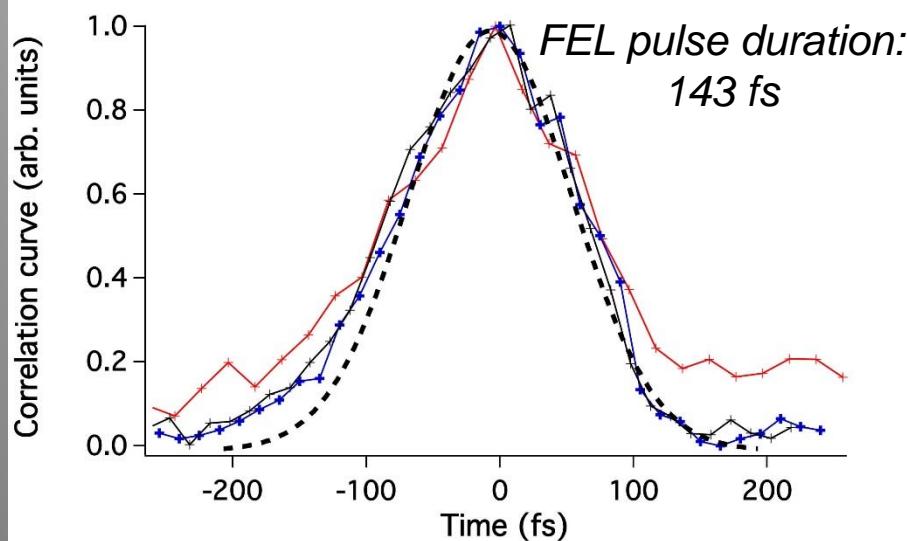
FEL pulse duration: 143 fs

Compressed FEL pulse

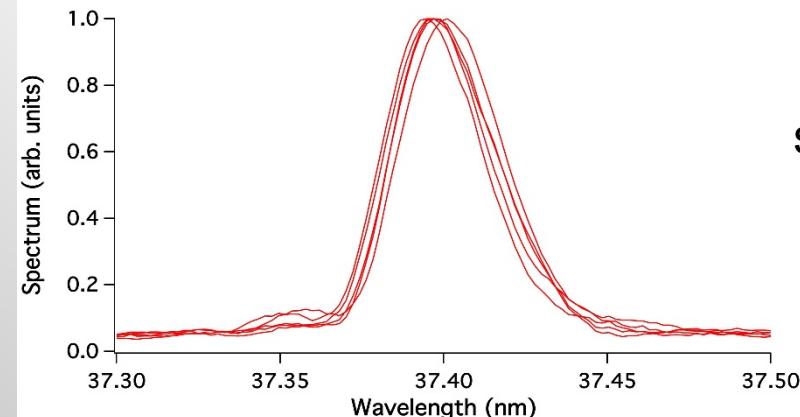


no compression

compression



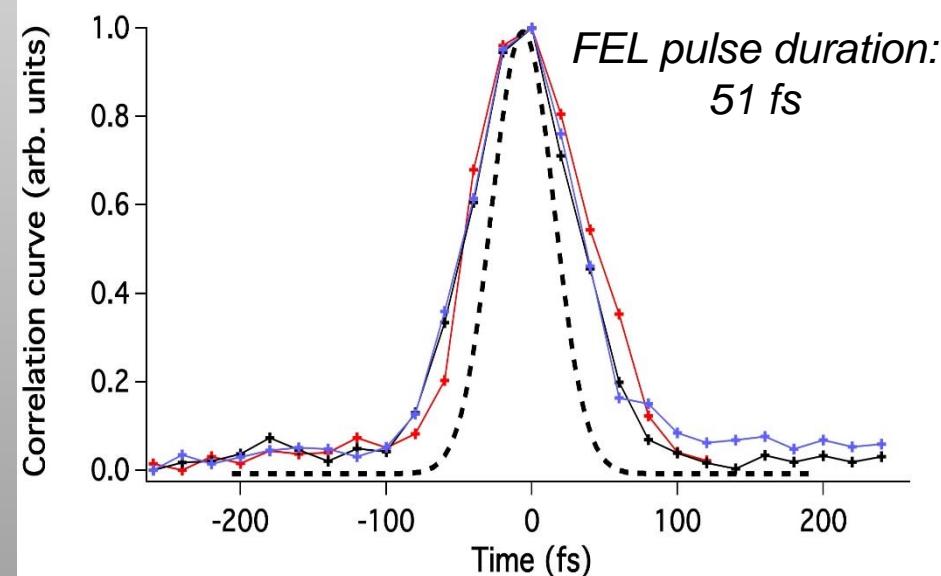
Compressed FEL pulse



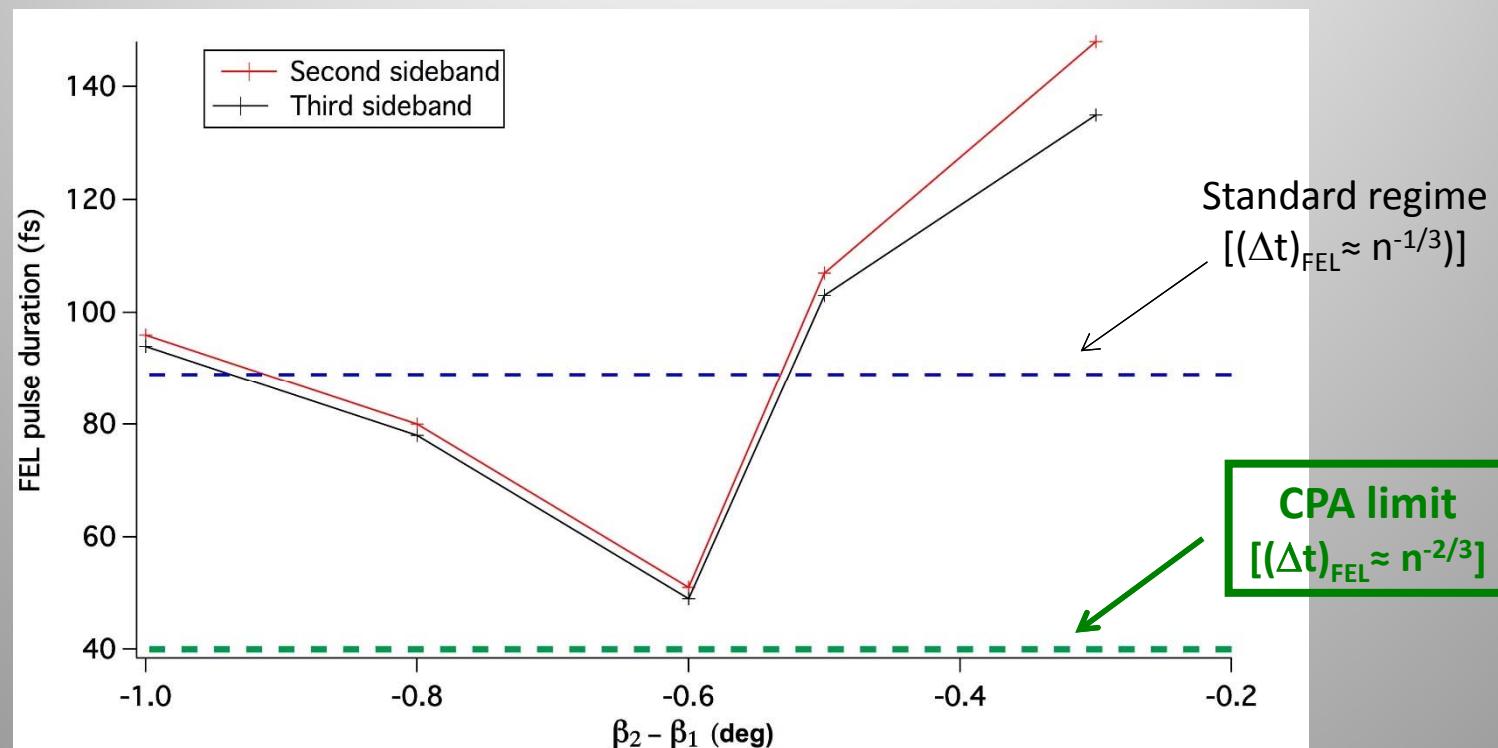
spectral width: 0.0446 nm

Time-bandwidth product
factor 1.1 above transform limit

compression



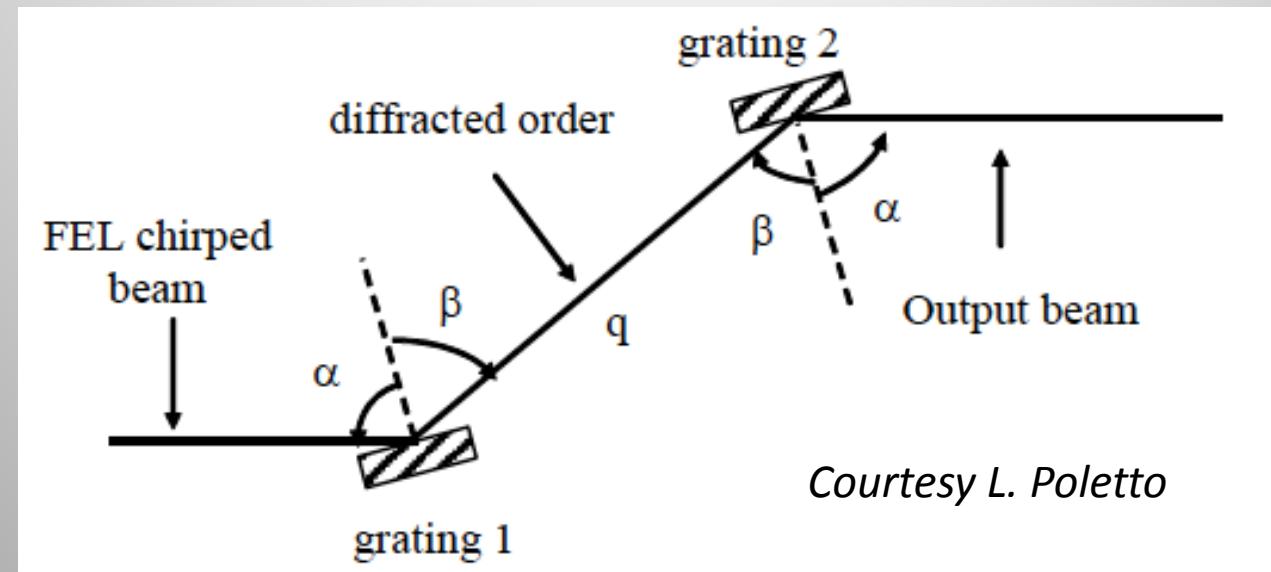
Measured FEL pulse duration as a function of the difference between the incidence angles on the two gratings.



The result obtained at maximum compression is close to the theoretical limit

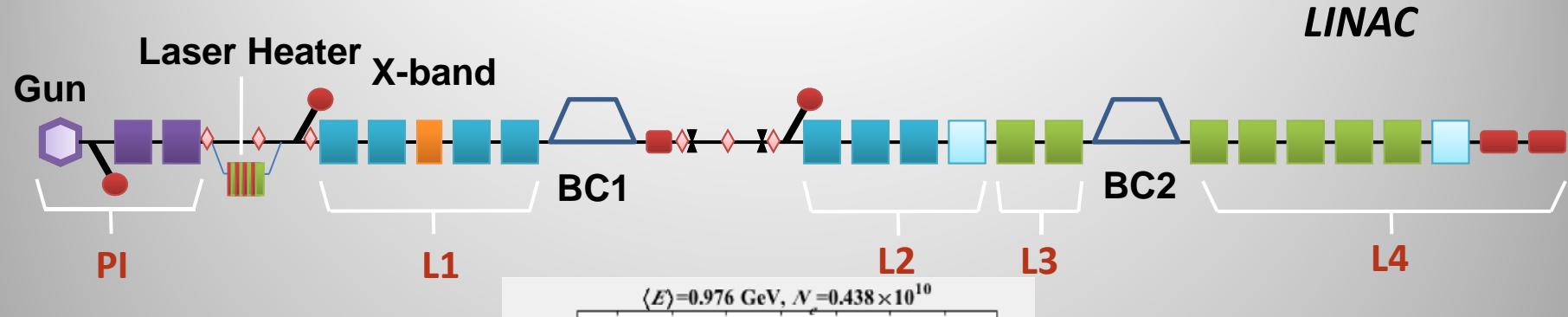
D. Gauthier et al., [Nature Communications](#),
vol. 7, (2016) p. 13688, doi :10.1038/ncomms13688(2016).

Next Step : CPA with FEL-2 @17 nm



The compressor can be used down to 4 nm with different gratings and with minor changes to the mechanics.

The constraints on the available space at FERMI force the use of the gratings in the classical geometry, which has intrinsically a low efficiency (in the range of few %).

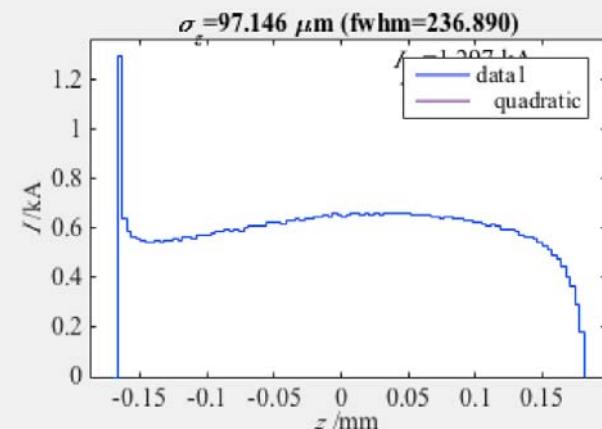
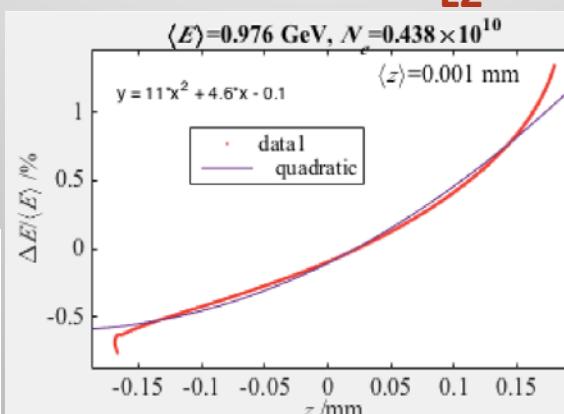


SETTINGS

- $Q=700\text{pC}$
- $L01: -28.5\text{deg}$
- $X: 15\text{MeV}$
- $BC1: R56=-41\text{mm}$
- $L02: \text{on-crest}$
- $L03: -45\text{deg}$
- $L04: k10-12 \text{ at } -50\text{deg}, k13-k15 -90\text{deg}$

Lin chirp: 13.5 MeV/ps

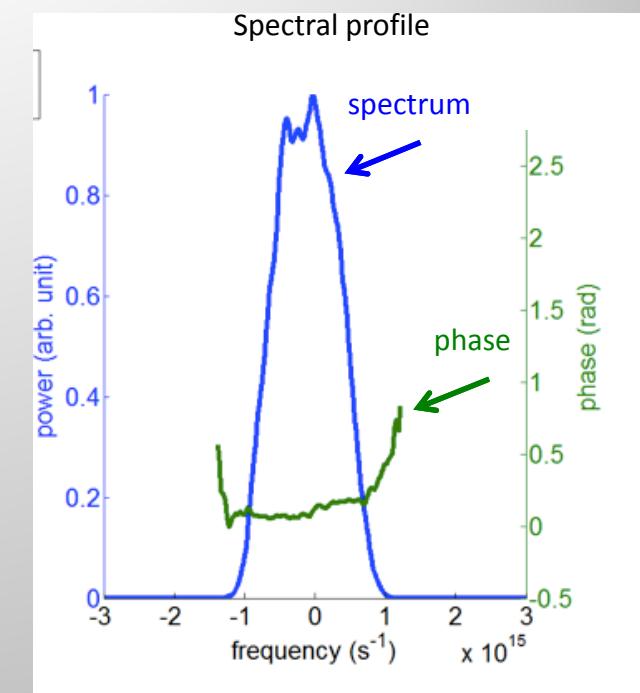
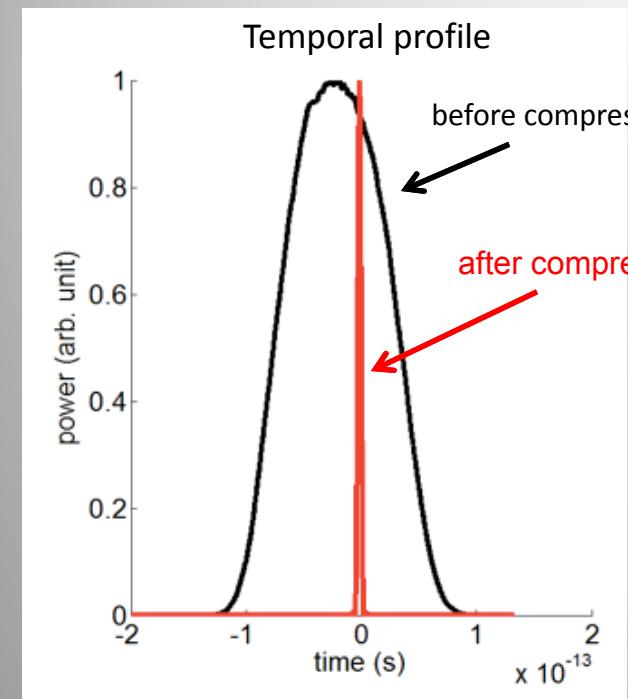
Quadr chirp: 9.7 MeV/ps^2



Courtesy G. Penco(numerical code: LiTrack)

For the design of the experiment, we used both seed and electron-beam parameters already available at FERMI.

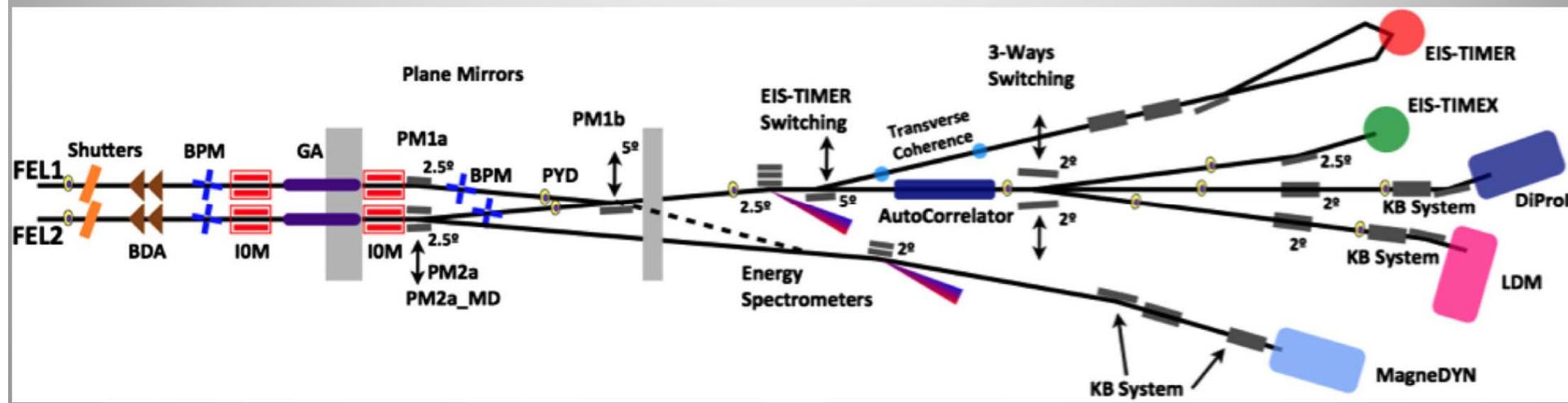
Beam and Laser Main Parameters			
Harmonic Number (n)	FEL Wavelength	15 (5x3)	17.6 nm
Seed Bandwidth	Pulse Duration (Fourier Limited)	1.8 nm	55fs (FWHM)
Duration Streched Seed	Chirp rate	300 fs (FWHM)	
Seed Power	Laser Energy in the Modulator	450 MW	150 μJ
Electron Beam Energy	Beam Current	0.98 GeV	600 A
Electrons Linear Chirp	Electrons quadratic Chirp	13 MeV/ps	10 MeV/ps²



Courtesy D. Gauthier (numerical code: Perseo)

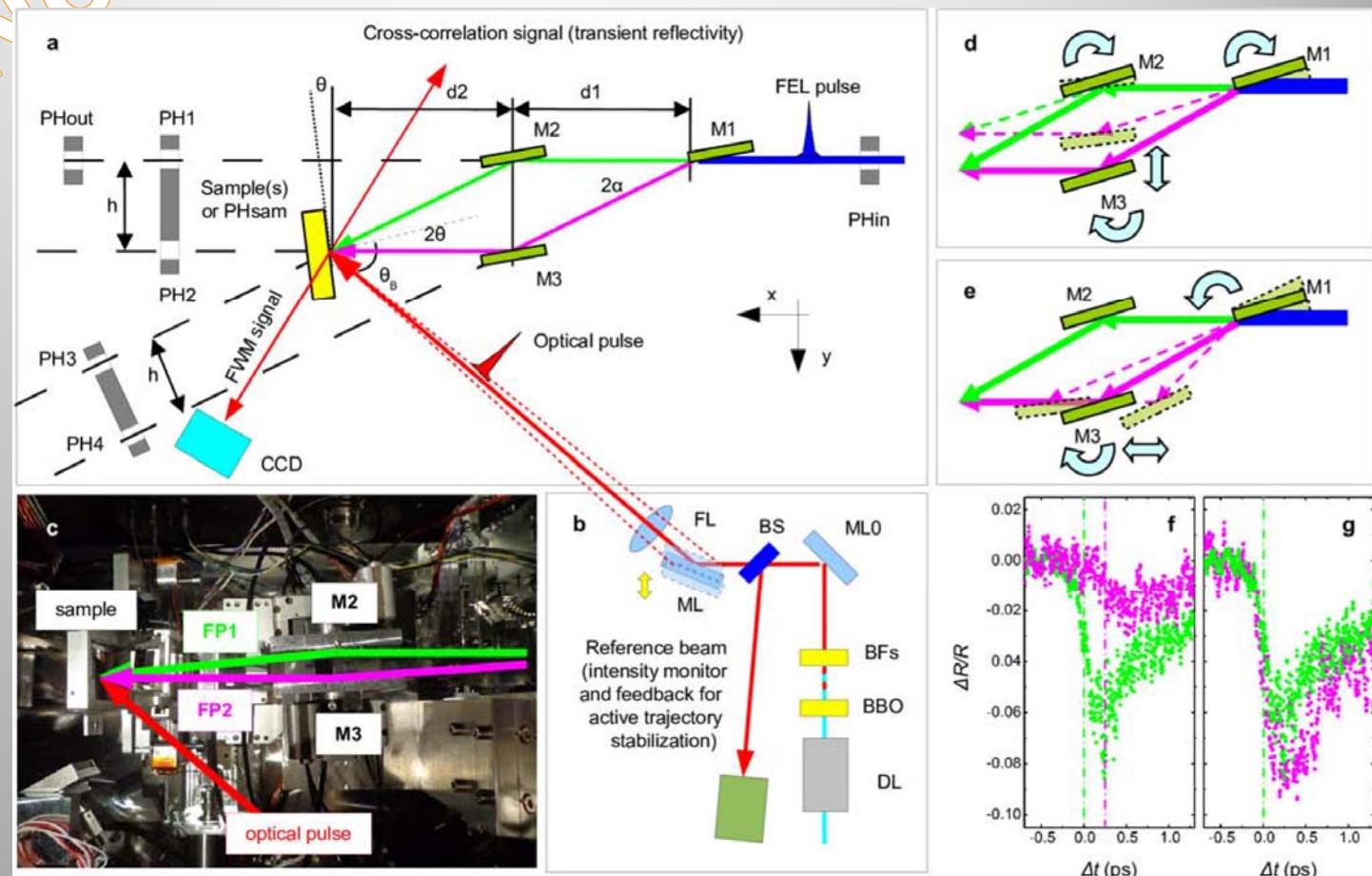
Expected FEL Performances			
FEL Output Bandwidth	FEL Output Duration	$2.5 \times 10^{14} \text{ rad/s}$	100 fs
Compensated Group Delay Dispersion (β)			400 rad/fs ²
Duration after Compression			9 fs
Output FEL pulse Energy			50 μJ
Estimated Peak Power vs. compressor Efficiency		1% = 170 MW	5% = 850 MW

New!



New!

Transient Grating approach



F. Bencivenga et al., Nature , 520, April 2015, 205

- -**High-order terms in the chirp** carried by the seed and/or in the electron beam energy distribution.
- Systematic errors, e.g., due to electron-beam phase-space curvature and/or non-homogeneous optics, may be compensated using a system (e.g., a deformable mirror) for active pulse shaping of the seed in a dispersive region.
- Stochastic errors, e.g., the ones due to the micro-bunching instability, will be more difficult to compensate. In order to mitigate the problem, we are working on the improvement of the quality of the photo-injector laser and will consider the possibility of operating the FEL at lower current.
- For extreme compressions, also the **jitter between the electrons and the seed** may play a role.

We demonstrated the possibility to carry out chirped pulse amplification in seeded FELs.

The technique, which can be extended to FELs based on self-amplified spontaneous emission, allowed us to achieve an unprecedented reduction of the FEL pulse duration, with respect to the input seed.

The relatively large seed pulse duration, as well as the quite low transmission efficiency of the compressor, prevented the generation, after compression, of very powerful pulses. However, by shortening the seed pulse and improving the compressor efficiency (e.g., by adopting an off-plane mount geometry of the gratings), the potential is there to produce, with existing technology, coherent few-fs/sub-fs gigawatt laser pulses inside the water window.

As an example, starting with a 10-fs seed pulse at 260 nm and tuning the FEL radiator at 4 nm ($n=65$), CPA would allow to generate a few-GW pulse of about 600 as.

The FERMI CPA collaboration includes:

- ✓ the following members of the FERMI team: E. Allaria, C. Callegari, C. Coreno, I. Cudin, M.B. Danailov, A. Demidovich, G. De Ninno, S. Di Mitri, B. Diviacco, E. Ferrari, P. Finetti, D. Gauthier, L. Giannessi, N. Mahne, G. Penco, L. Raimondi, P. Rebernik, R. Richter, E. Roussel, L. Sturari, C. Svetina, M. Trovò, M. Zangrando
- ✓ The CNR of Padova (Italy): F. Frassetto, P. Miotti, L. Poletto
- ✓ The LOA of Palaiseau (France): H. Dacasa, B. Mahieu, P. Zeitoun
- ✓ The CEA of Saclay (France): D. Garzella
- ✓ The IPFN of Lisbon (Portugal): M. Fajardo, S. Kunzel
- ✓ The European XFEL (Germany): M. Meyer, T. Mazza
- ✓  Connection in progress with The CNRS-ISMO : A. Klisnick, A. Le Marec